

# *Quantum Error Correction in correlated quantum noise*

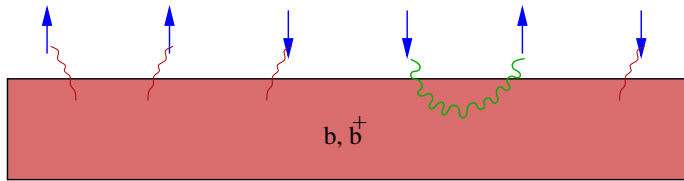
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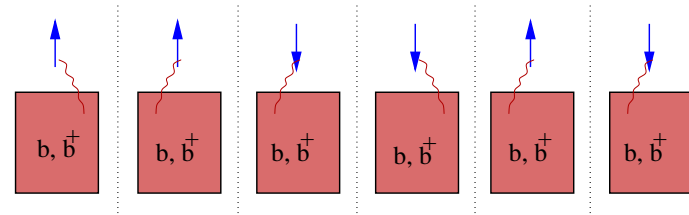
*In collaboration with:*

- Sandra Frank ( → FZ Jülich)
- Stefan Borghoff
- Thomas Zell

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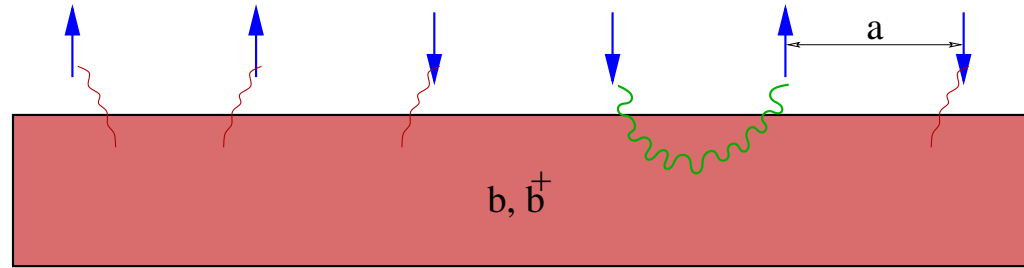
Outline:

*I. Physical model*

*II. Effect of noise correlations on QEC: leading-order estimate*

*III. Results of exact calculation*

I. Physical model :



- $n$  (effective) spin- $\frac{1}{2}$
- common bosonic bath
- linear and local spin-boson coupling :
- no direct coupling between spins

$$\sum_l g S_l \otimes A_l$$

$$S_l = X_l, Y_l \text{ or } Z_l$$

$A_l =$  bosonic field strength at spin  $l$

noise correlations due to *boson exchange* !

QEC :

- CSS codes for  $\leq t$ -qubit errors
- instantaneous
- ideal

## To which extent do noise correlations interfere with QEC ?

II. Leading-order estimate :

$$\psi \xrightarrow{\mathcal{N}_\tau} \psi' \xrightarrow{\mathcal{R}} \psi''$$

$\mathcal{N}_\tau$  : noise on  $n$  qubit register during time  $\tau$

$\mathcal{R}$  :  $\leq t$ -qubit error correcting operation

residual error :

$$\Delta \equiv \|\psi'' - \psi\| \sim \binom{n}{t+1} (\alpha g^2)^{t+1} \langle A_{l_1}^2 \dots A_{l_{t+1}}^2 \rangle,$$

$g$  : coupling constant

$\alpha = \alpha(\tau)$

$$\Delta \sim \binom{n}{t+1} (\alpha g^2)^{t+1} \langle A_{l_1}^2 \dots A_{l_{t+1}}^2 \rangle,$$

$t$  : number of correctable errors

*independent noise* :

$$\langle A_{l_1}^2 \dots A_{l_{t+1}}^2 \rangle = \langle A_{l_1}^2 \rangle \dots \langle A_{l_{t+1}}^2 \rangle = \langle A^2 \rangle^{t+1}.$$

infinite spin distance

*correlated noise* :

$$\langle A_{l_1}^2 \dots A_{l_{t+1}}^2 \rangle = \langle A^{2(t+1)} \rangle = \underline{(2t+1)!!} \langle A^2 \rangle^{t+1}$$

vanishing spin distance

$$\implies \frac{\Delta_{corr}}{\Delta_{ind}} = (2t+1)!!$$

→  $n, t \gg 1$ , constant correctable error rate  $t/n = q$  :

*independent noise:*

$$\Delta \sim \binom{n}{q n} \left( \alpha g^2 \langle A^2 \rangle \right)^{qn} \approx \left( c_q \alpha g^2 \langle A^2 \rangle \right)^{qn} \xrightarrow{n \rightarrow \infty} 0$$

if  $g^2 < g_0^2 \equiv (c_q \alpha \langle A^2 \rangle)^{-1}$

*correlated noise:*

$$\Delta \sim \binom{n}{q n} (2qn + 1)!! \left( |g|^2 \langle A^2 \rangle \right)^{qn} \approx \left( c_q \alpha n g^2 \langle A^2 \rangle \right)^{qn} \xrightarrow{n \rightarrow \infty} \infty$$

for all  $g^2 > 0$  !

$$\binom{n}{q n} \approx e^{nH_2(q)}, \quad (2n + 1)!! \sim (2n)^n$$

## Conclusion:

- correlations due to boson exchange do matter
- perturbative series in  $g^2$  problematic
- unboundness of  $A^2$  essentially

$$\frac{\Delta_{corr}}{\Delta_{ind}} \simeq \frac{\langle A^{2t} \rangle}{\langle A^2 \rangle^t} \simeq (2t)^t$$

$$\lim_{n \rightarrow \infty} \Delta_{corr} = \infty !?$$

$$\text{otherwise } \lim_{t \rightarrow \infty} \frac{\langle A^{2t} \rangle}{\langle A^2 \rangle^t} = 1$$

→ *investigation of exactly solvable model: III.*

### III. Dissipationless $n$ -spin-boson model:

$$H = \frac{\varepsilon}{2} \sum_{l=1}^n Z_l + \sum_k \omega_k b_k^\dagger b_k + \sum_{k,l} Z_l (g_k e^{ik r_l} b_k^\dagger + h.c.)$$

Unruh '95; Palma, Suominen, Ekert '96

↪ exact noise operation  $\mathcal{N}_\tau$  on  $n$ -qubit register:  $\rho \xrightarrow{\tau} \mathcal{N}_\tau(\rho)$

#### Ensemble of all $[[n, k]]$ CSS codes:

- $n$  large
- fixed information rate  $R = k/n$
- correctable error rate  $q = t/n$  determined by  $R = 1 - 2H_2(2q)$

↪ ensemble averaged residual error :

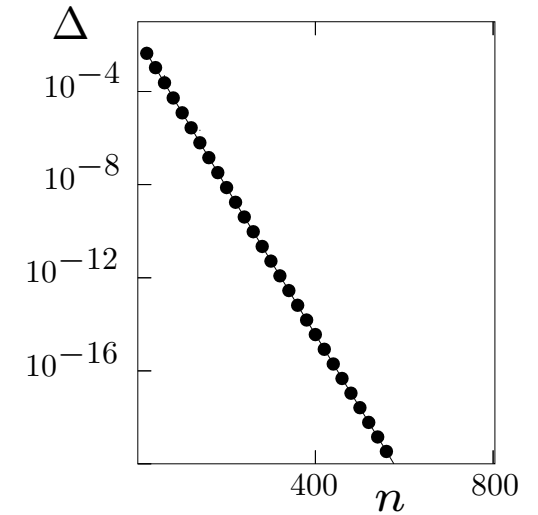
$$\Delta_{n,q} = 1 - \left\langle \int_C d\psi F(\psi, \mathcal{R}_C \circ \mathcal{N}(\psi)) \right\rangle_{n,q}$$



Independent noise ( $a = \infty$ ):

$$\Delta_{n,q} \simeq \exp \left[ -n \frac{(q - K(0,\tau))^2}{2K(0,\tau)} \right],$$

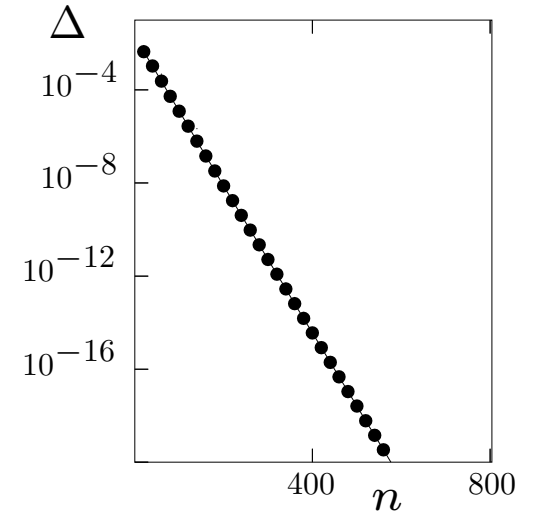
$K(r, \tau)$  : time- and distance-dependent decoherence coefficient



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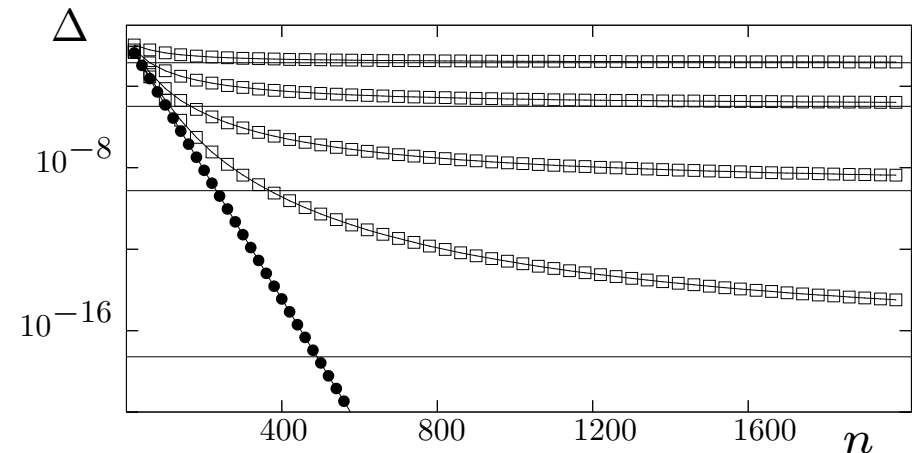
$K(r, \tau)$  : time- and distance-dependent decoherence coefficient



Correlated noise ( $a < \infty, n \gg 1$ ):

$$\Delta_{n,q} \sim \exp \left[ \frac{-q}{K(r, \tau)} \right]$$

$r$  : spatial extension of register



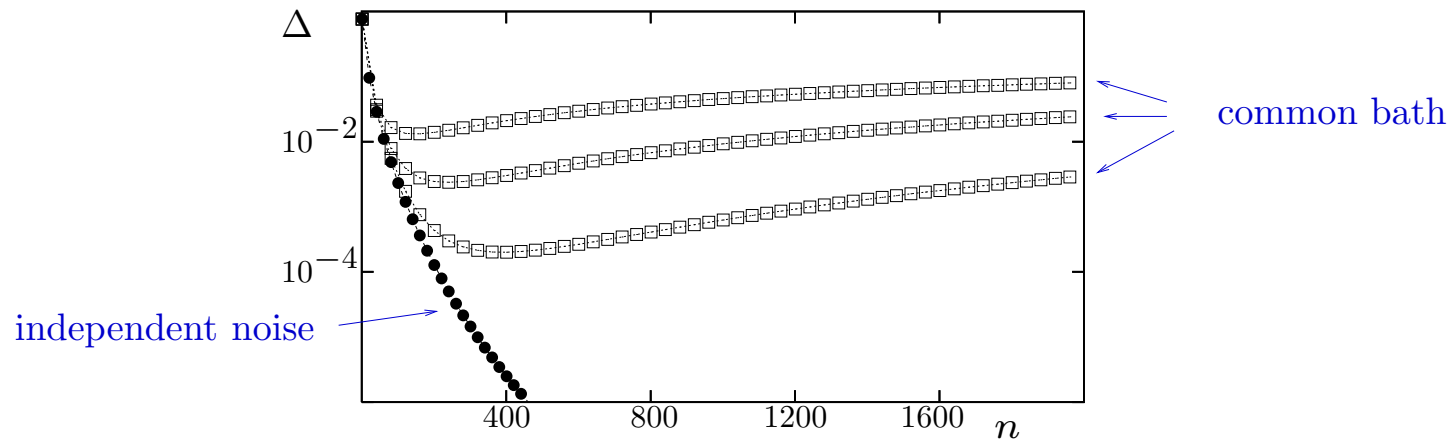
$q = 0.05$ , fixed  $K_{r,\tau} = 0.01, 0.005, 0.0025$ , and  $0.00125$ ,  $K_O = 0.01$

Example:

- $n$  spins on cubic array
- spatial extension  $r = n^{1/3} r_0$
- observation time  $\tau = n^{1/3} \tau_0$
- ohmic bosonic bath:

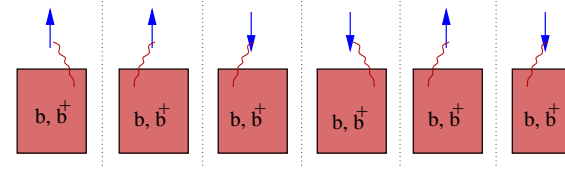
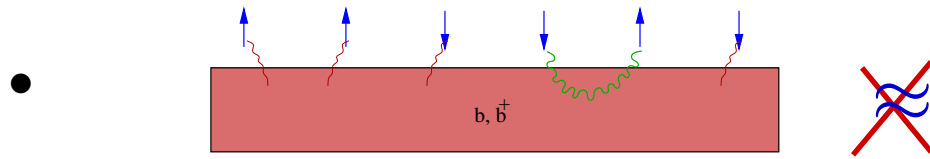
(since  $\tau \propto r$  !?)

$$K(n) \equiv K(n^{1/3} r_0, n^{1/3} \tau_0, T) \stackrel{!}{=} n^{1/3} K(r_0, \tau_0, T)$$



CSS error correction significantly reduced by boson-exchange correlations

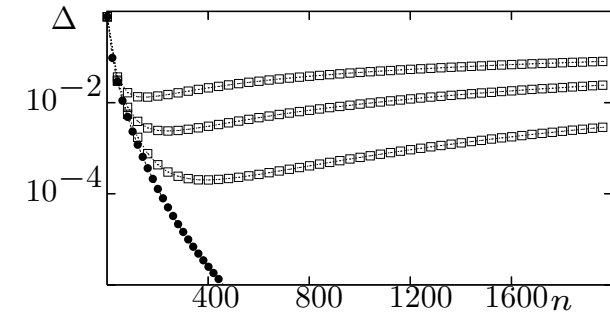
## Summary/Conclusion



$$\langle A^{2t} \rangle \gg \langle A^2 \rangle^t$$

- spin-boson model + CSS error correction:

$$\Delta = 1 - \langle F_{av}(C, \mathcal{R} \circ \mathcal{N}) \rangle_{[[n, qn]]}$$



## Open questions

- improvement by use of “decoherence-free-subspaces” ?
- impact on FT-QC?
- relation to improved Threshold Theorems?

cf. B. Terhal *et al.* '05, P. Aliferis *et al.* '05, D. Aharonov *et al.* '06

Lit.: RK, S. Frank, PRL **95**, 230503 '05  
 RK, PRA **75**, 062315 '07  
 S. Borghoff, RK, PRA **76** '07